

Flame Acceleration, DDT, and Hydrogen in Channels

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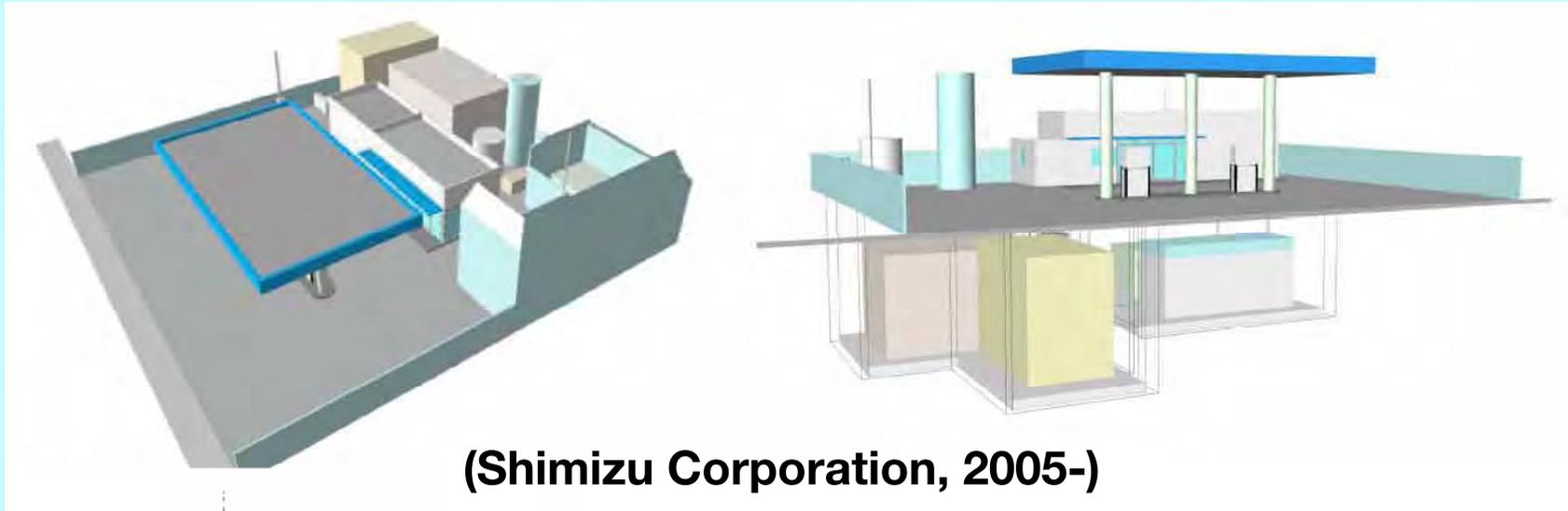
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Geraint O. Thomas (Leeds University, *Abershock*)

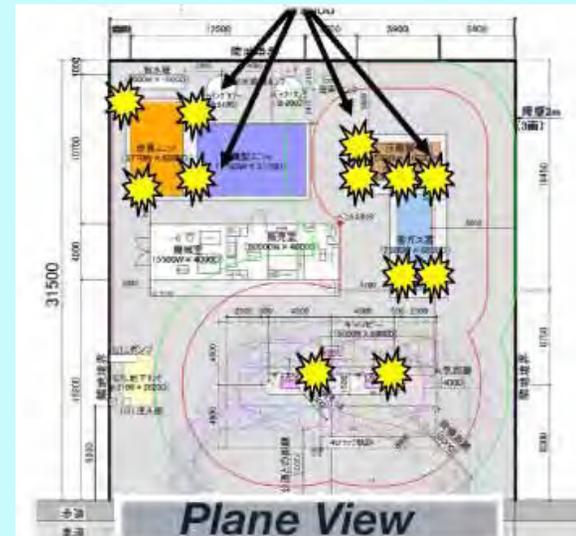
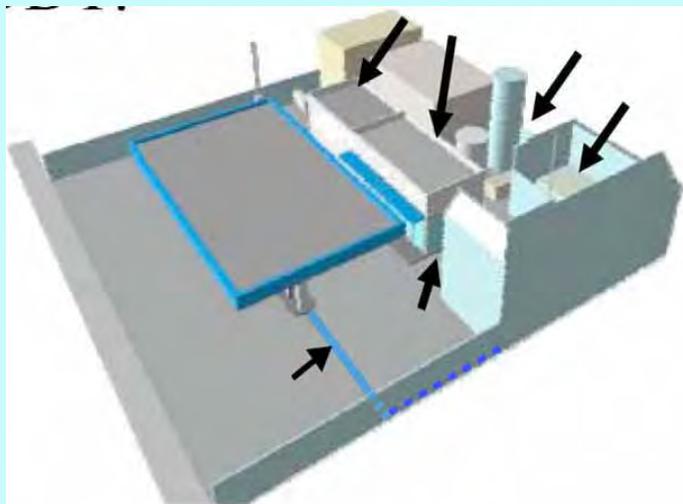
Sponsored by:

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Proposed Hydrogen Fuel Station

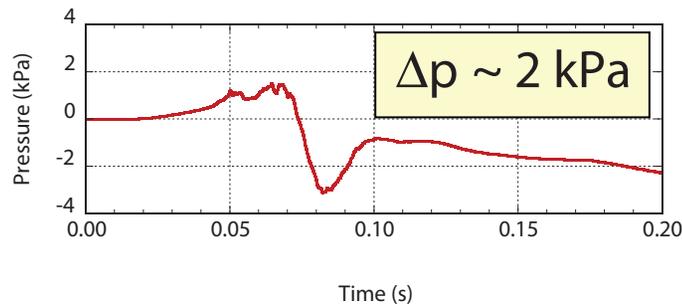


Fuel leakages could lead to flames and even detonations ...

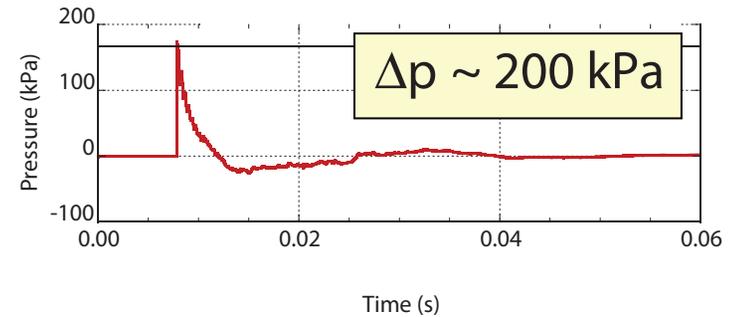


Deflagration and Detonation

Deflagration



Detonation



Pressure histories at 4m away from the ignition point for 5.27m³ of a H₂-air mixture



- Window glass shatters



- Concrete shatters
- Brick shears apart
- Lung damage
- Eardrum rupture

A flame ignited in leaked H₂ gas could undergo a deflagration-to-detonation transition !

Focus of Recent Research

Answer the question:

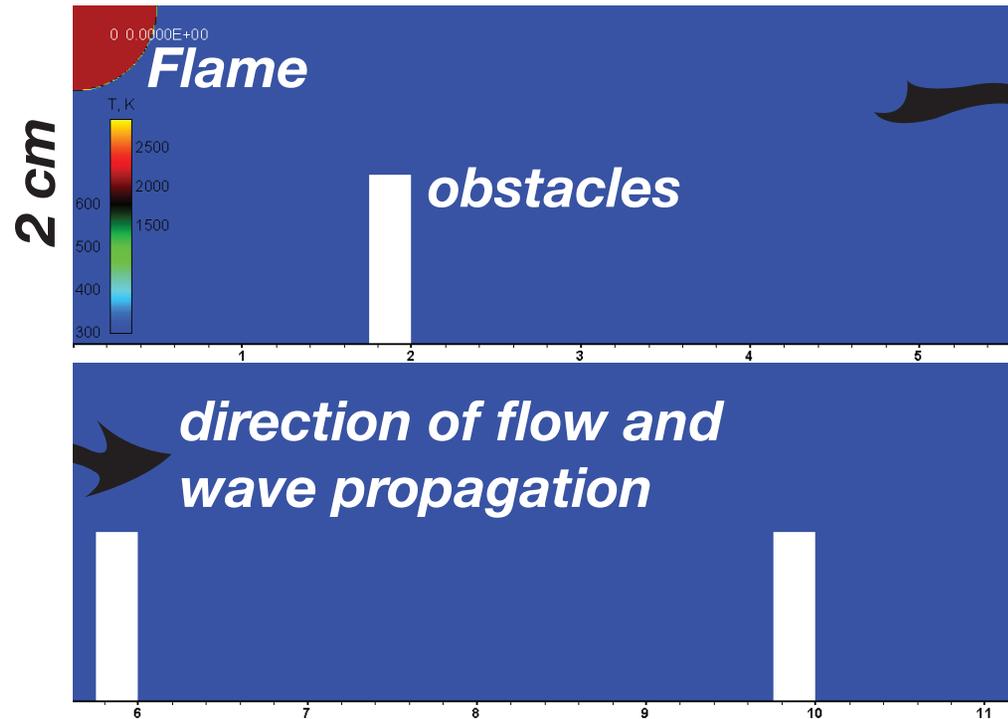
How confidently ... qualitatively or quantitatively ... can we compute hydrogen combustion from ignition, to flame acceleration, and even to DDT?

This Involves:

- Developing and test low-cost, low-order chemical-diffusive models**
- Developing and test high-order methods with AMR**
- Testing performance of multidimensional, reactive numerical models**

H₂-Air Mixture Ignited in a Channel with Obstacles

Beginning of Movie:

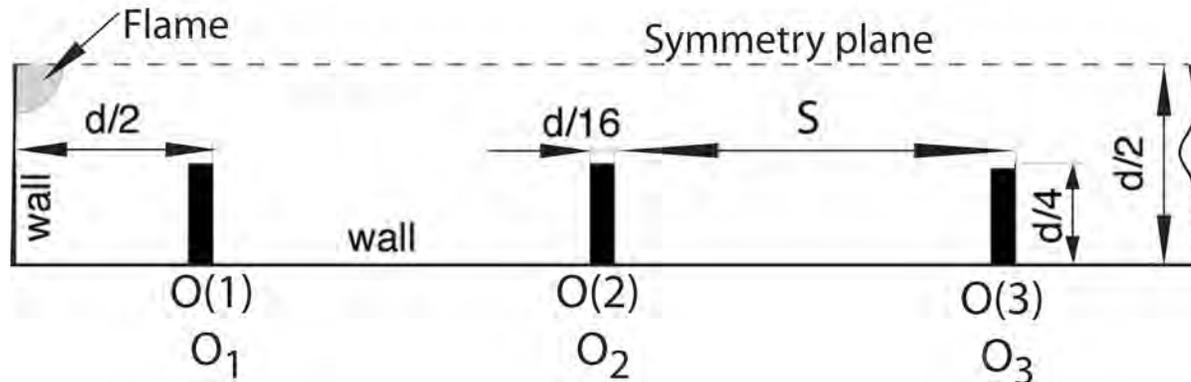


Movie will show how ...

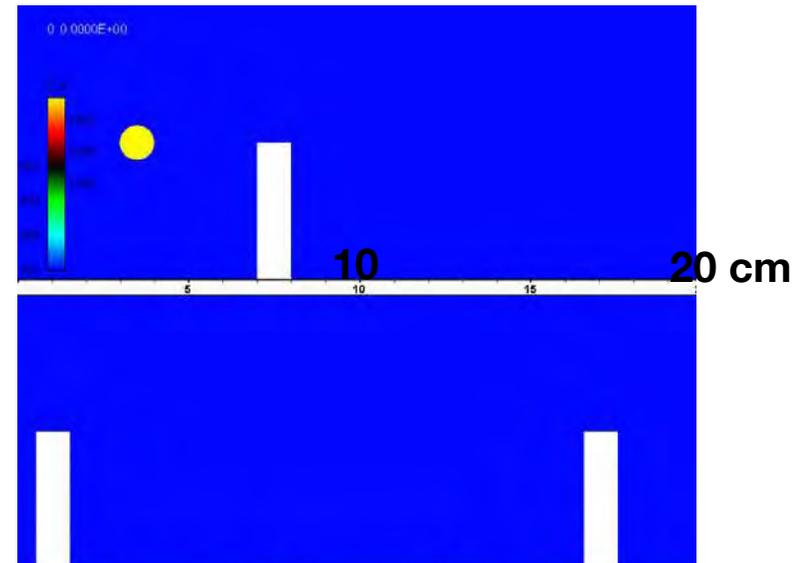
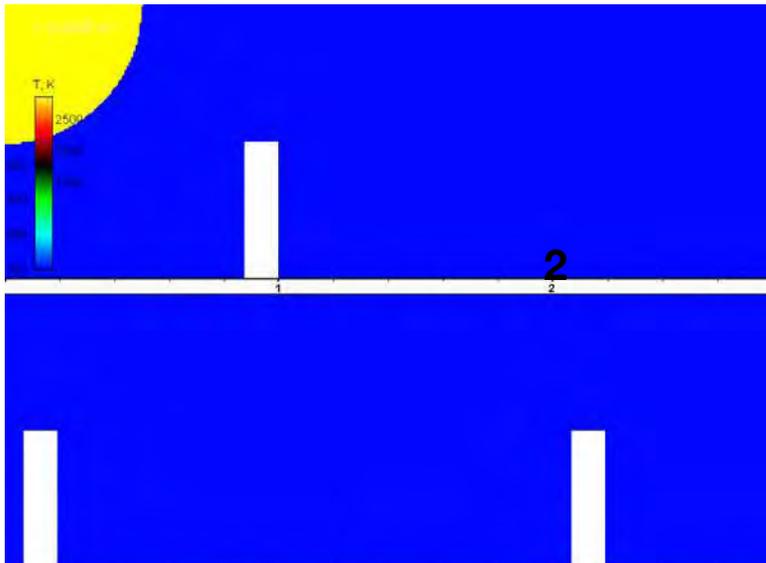
... starting with spark ignition in enclosures containing combustible mixtures, turbulent flames develop and produce shock waves. This leads to the formation of unsteady shock-flame complexes and detonations.

Channels with Obstacles

Channels filled with stoichiometric H₂-Air



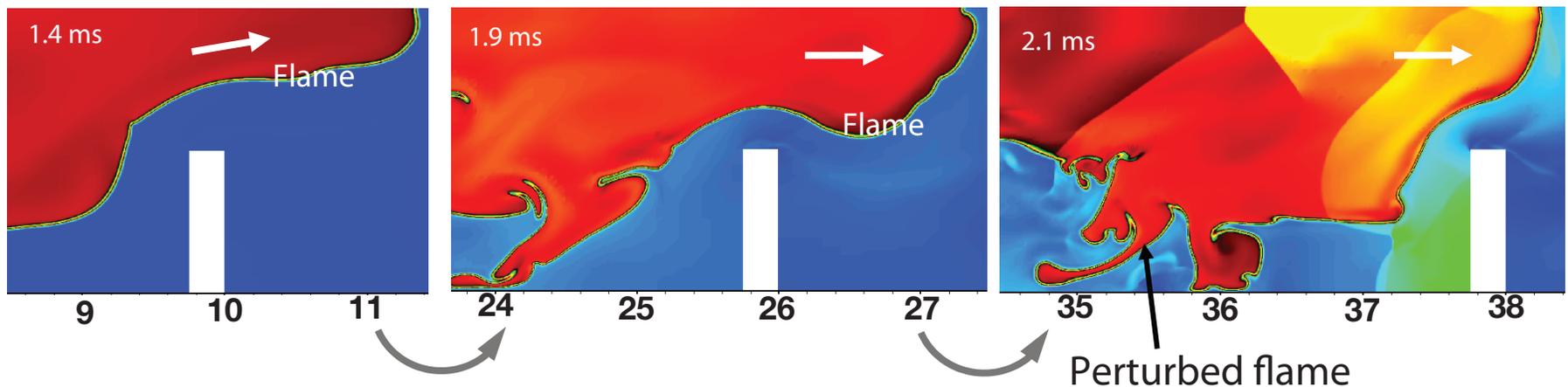
$d/2 = 1, 2, 4, \text{ or } 8 \text{ cm}$; **Constant blockage ratio**
Smooth or spark ignition; Constant input energy



Computational Result

Early Flame Propagation

Temperature Contours

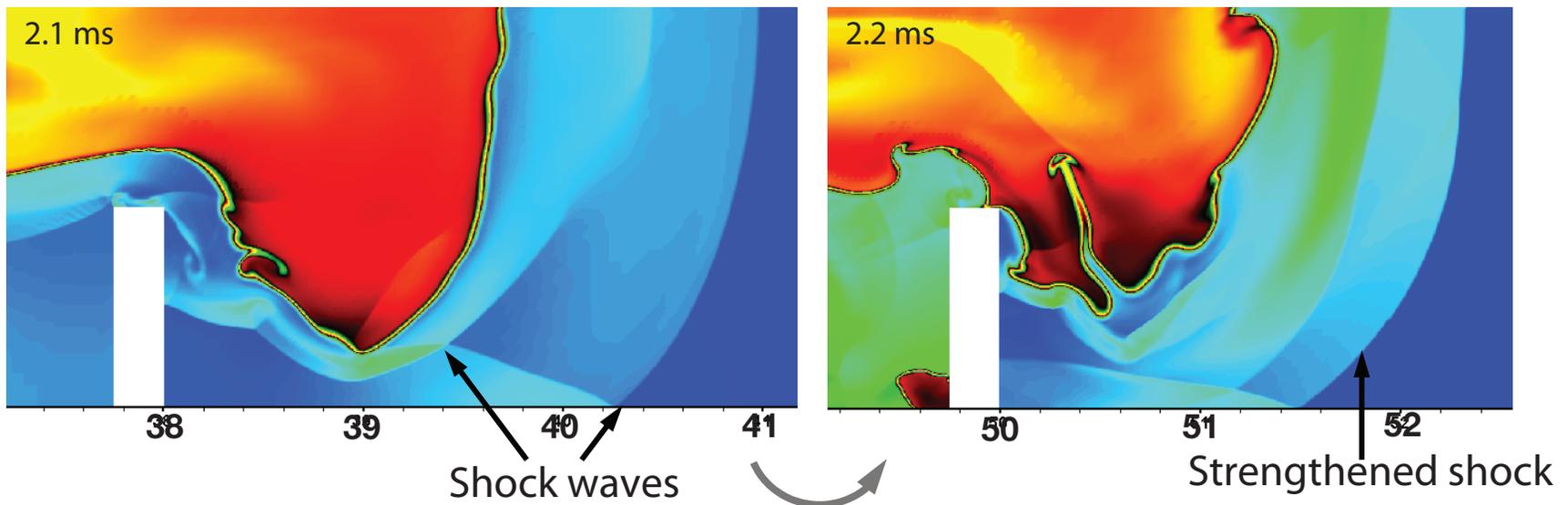


- The initially laminar flame moves slowly into the unreacted material (to the right).
- Obstacles perturb the flow, which then interacts with and distort the flame, so that the flame becomes turbulent.

Computational Result

Shock Wave Formation

Temperature Contours

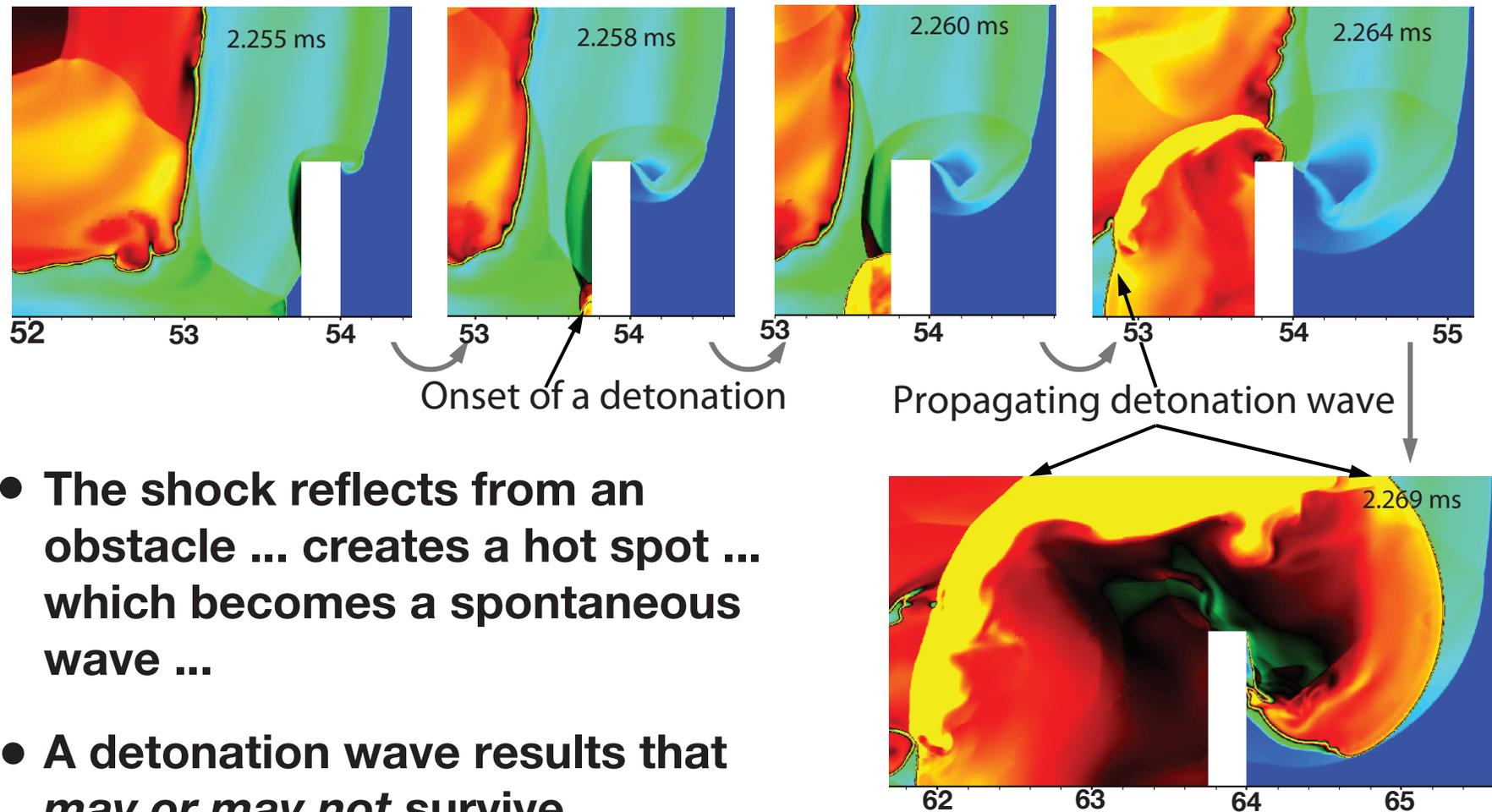


- The turbulent flame generates compression waves, which eventually coalesce to form a shock in front of the flame.
- The shock is continuously strengthened by compression waves coming from behind.

Computational Result

Transition to Detonation

Temperature Contours



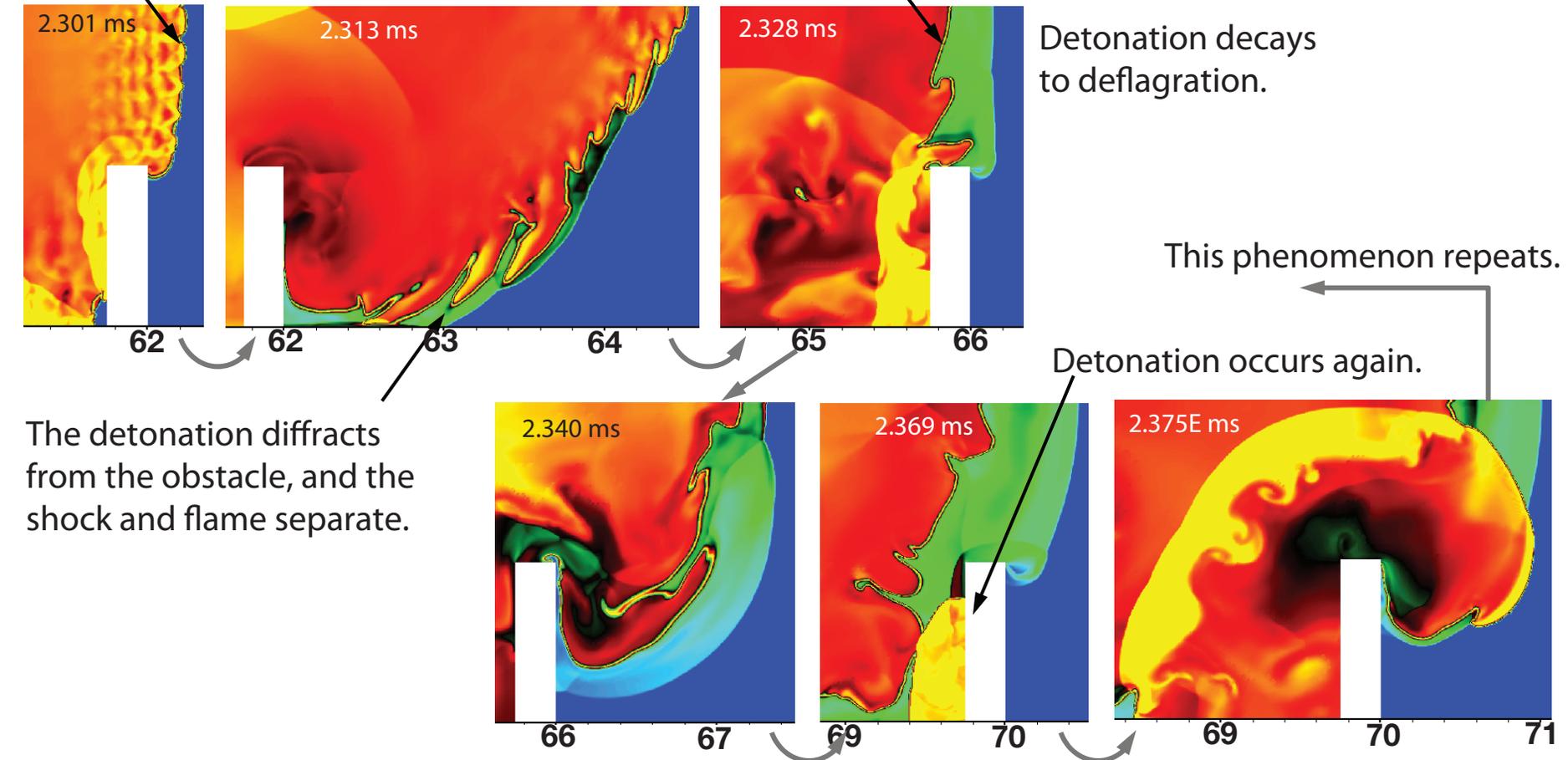
Computational Result

Detonation Wave Propagation

Quasi-detonation for $d = 4$ and 8 cm

Detonation wave

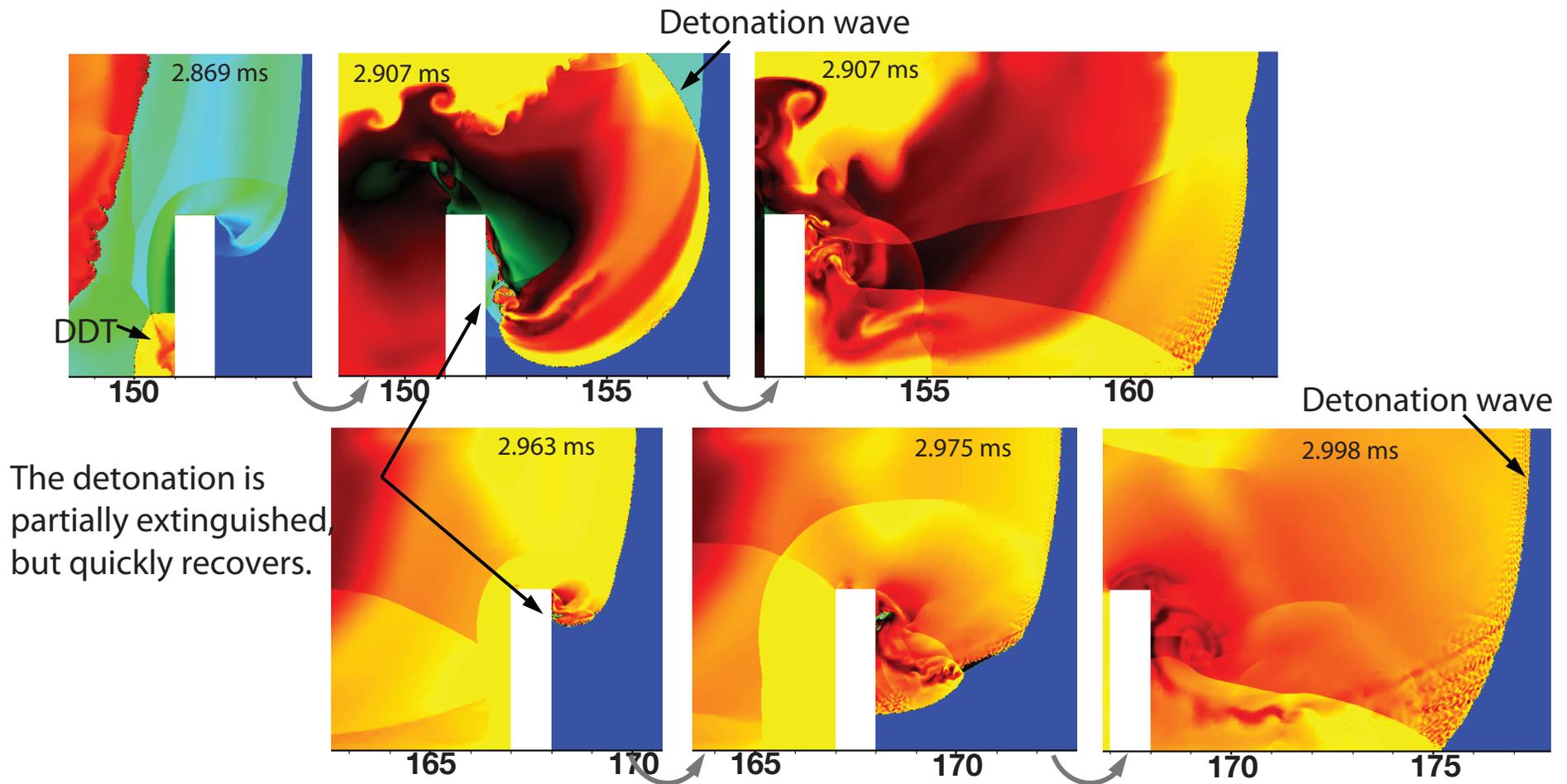
Flame



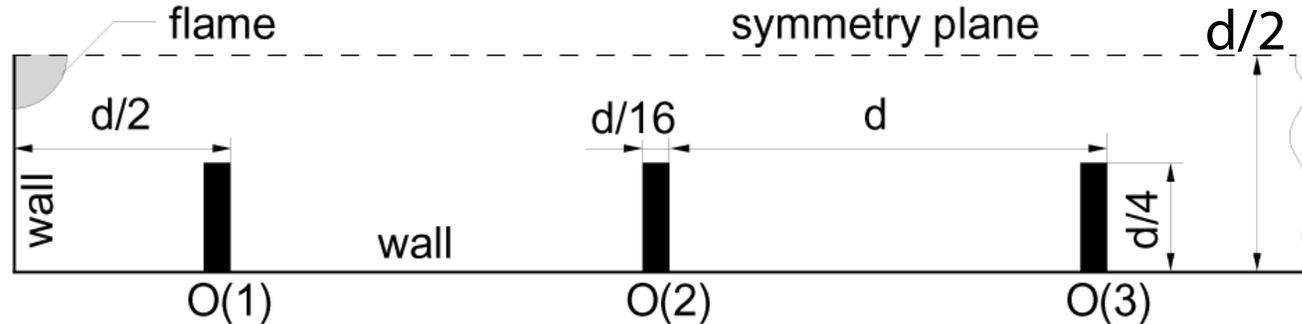
Computational Result

Detonation Wave Propagation

For large enough channels, the detonation successfully propagates over the obstacle.



Regimes of Flame Propagation



1. Slow deflagrations

subsonic

2. Fast deflagrations (“choking”)

$1/3 - 1/2 D_{CJ}$,
 $\leq c_s$ in burned gas

Flame decoupled from leading shock, spreads through molecular or turbulent diffusion and convection.

3. Quasi-detonations

$< D_{CJ}$

Shock & flame coupled some sometimes when reaction triggered by shock compression. Sometimes, shock and flame decoupled as detonation diffracts over obstacles. Observed when width of unobstructed part of channel $>$ few times l . Propagation velocity increases with d/l , and reaches D_{CJ} when propagation becomes independent of diffraction effects.

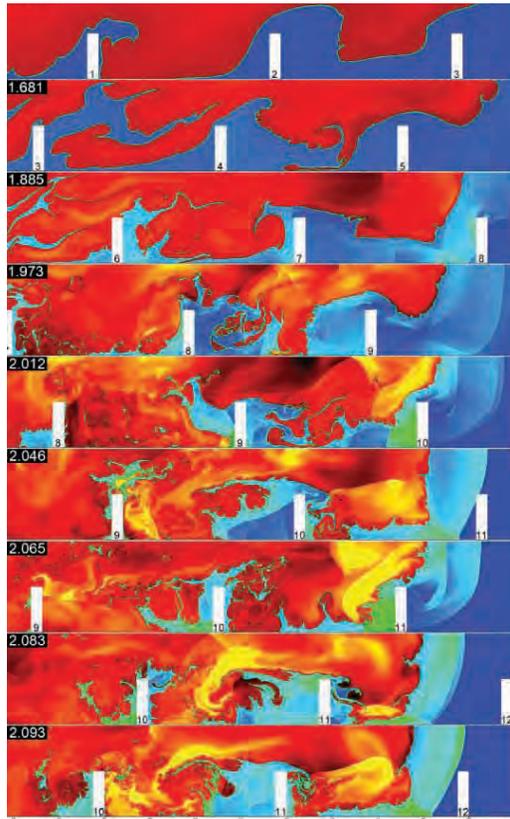
4. Detonations

D_{CJ}

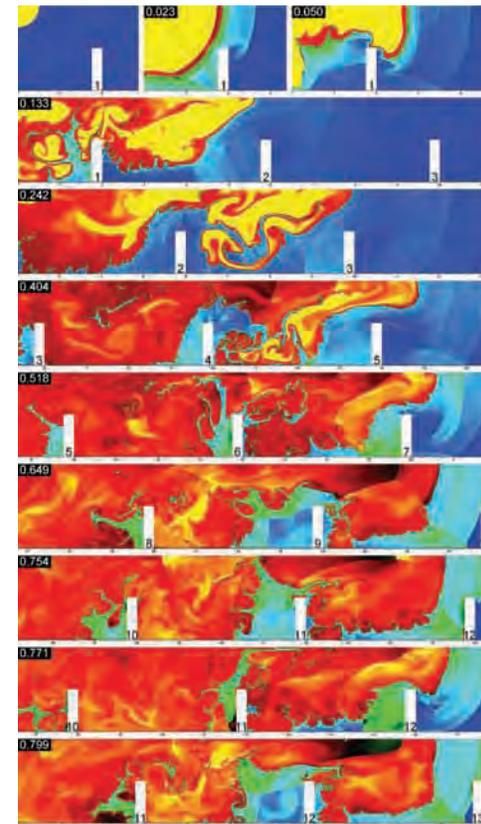
H₂-Air Mixture Ignited in a Channel with Obstacles

Two Modes of Ignition

Smooth Flame Ignition



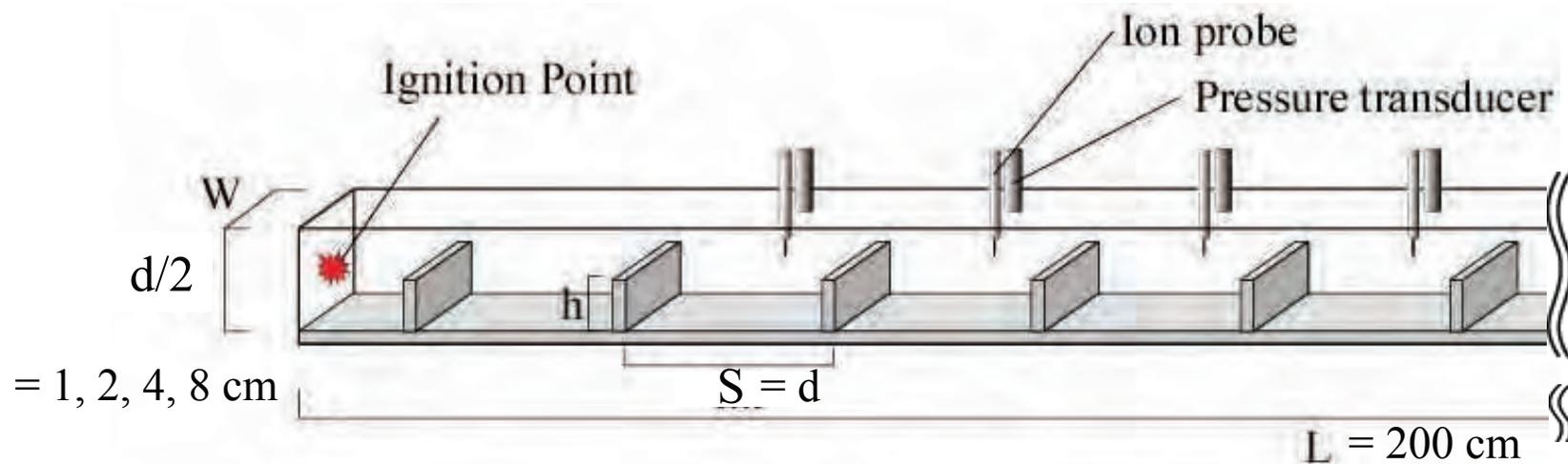
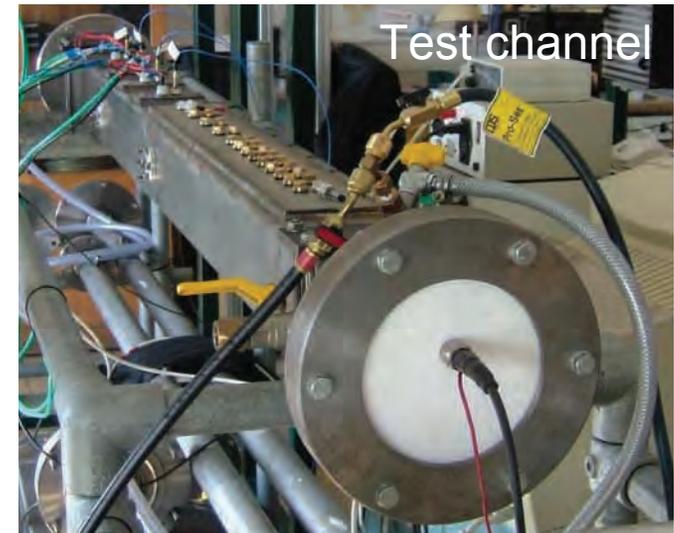
Spark Ignition



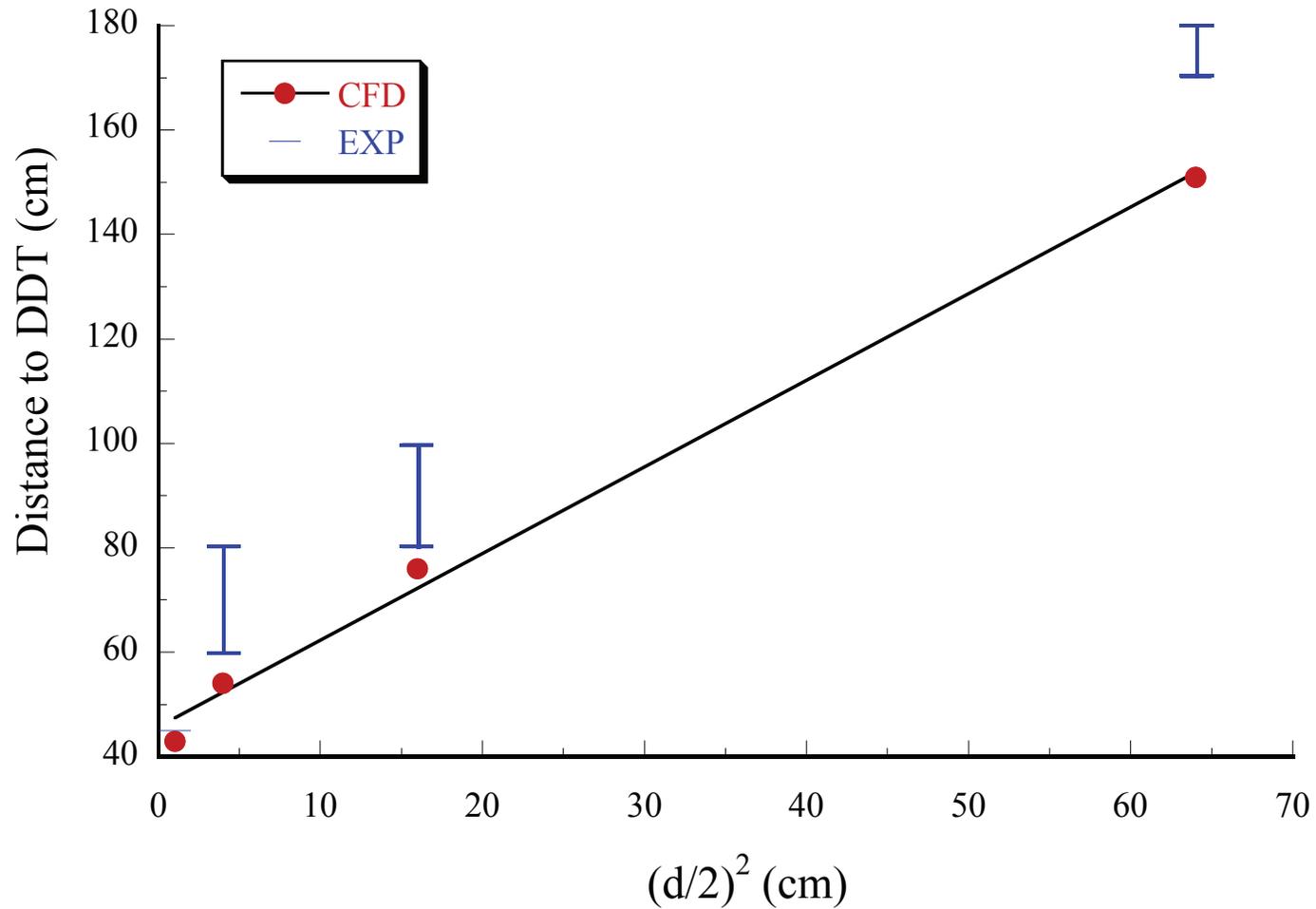
Initial development appears quite different: Immediate effects of shocks apparent for spark ignition.
Time to DDT less for spark ignition (0.61 ms from 0.76 ms)
Location is essentially the same.

Experimental Tests

- **DDT in a channel with obstacles**
 - For example, piping space
 - The same conditions as computations
- **Conducted by Andrej Teodorczyk
Warsaw University of Technology**

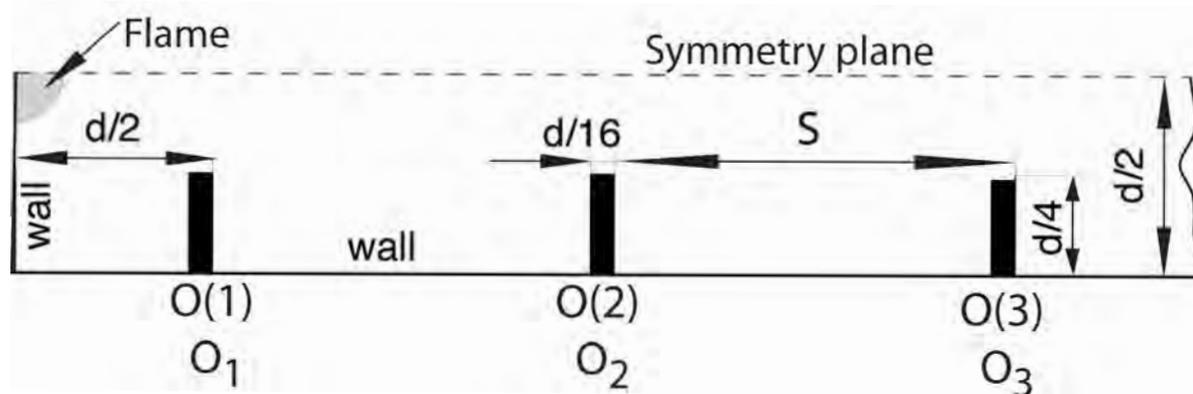


Effect of Channel Height on DDT Distance



The distance to DDT is proportional to d^2

Cases Considered for H₂-Air Channels



Variations on:

Form of ignition (smooth vs spark)

Channel height ($d/2$)

Blockage ratio

Obstacle spacing (S)

Obstacle symmetry (top and bottom)

Dimensionality (3D vs 2D)

Stoichiometry

Stochasticity

...

Back to the Basics

Chemical model

Effects of variable stoichiometry

Even more complex geometries

Why do simulations “work” ..?

Effects of non-Kolmogorov turbulence

Stochasticity

THE PHYSICAL MODEL

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{U}) = 0 \quad \text{c}$$

$$\frac{\partial (\rho \mathbf{U})}{\partial t} + \nabla \cdot (\rho \mathbf{U} \mathbf{U}) + \nabla P + \nabla \cdot \hat{\tau} = 0 \quad \text{c}$$

$$\frac{\partial E}{\partial t} + \nabla \cdot ((E + P) \mathbf{U}) + \nabla \cdot (\mathbf{U} \cdot \hat{\tau}) + \nabla \cdot (K \nabla T) = 0 \quad \text{c}$$

$$\frac{\partial (\rho Y)}{\partial t} + \nabla \cdot (\rho Y \mathbf{U}) + \nabla \cdot (\rho D \nabla Y) - \rho \dot{w} = 0 \quad \text{c}$$

$$\hat{\tau} = \rho \nu \left(\frac{2}{3} (\nabla \cdot \mathbf{U}) \hat{I} - (\nabla \mathbf{U}) - \nabla \mathbf{U} \right)^\dagger$$

$$P = \frac{\rho R T}{M} \quad \text{c} \quad E = \frac{P}{(\gamma - 1)} + \frac{\rho U^2}{2} \quad \text{c}$$

$$\frac{dY}{dt} \equiv \dot{w} = -A \rho Y \exp \left(-\frac{Q}{RT} \right)$$

$$\nu = \nu_0 \frac{T^n}{\rho} \quad \text{c} \quad D = D_0 \frac{T^n}{\rho} \quad \text{c} \quad \frac{K}{\rho C_p} = \kappa_0 \frac{T^n}{\rho}$$

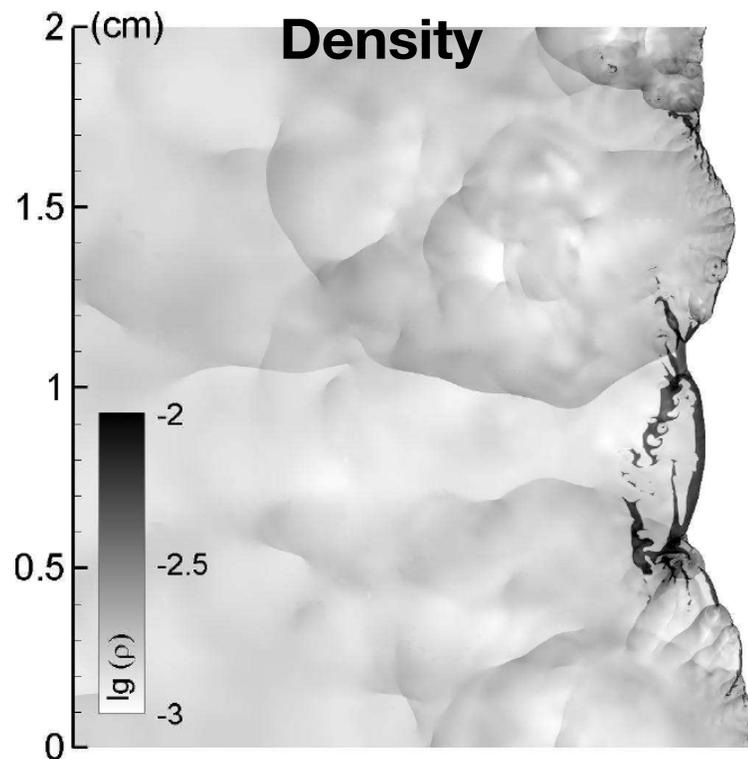
$$Le = \frac{K}{\rho C_p D} = \frac{\kappa_0}{D_0} \quad \text{c} \quad Pr = \frac{\rho C_p \nu}{K} = \frac{\nu_0}{\kappa_0} \quad \text{c} \quad Sc = \frac{\nu}{D} = \frac{\nu_0}{D_0}$$

Material, Chemistry, and Reaction Wave Parameters

Stoichiometric Hydrogen-Air

Quantity	Value	Definition
Input		
T_0	293 K	Initial temperature
P_0	1 atm	Initial pressure
ρ_0	$8.7345 \times 10^{-4} \text{ g/cm}^3$	Initial density
γ	1.17	Adiabatic index
M	21 g/mol	Molecular weight
A	$6.85 \times 10^{12} \text{ cm}^3/\text{g-s}$	Pre-exponential factor
$E_a (= Q)$	$46.37 RT_0$	Activation energy
q	$43.28 RT_0/M$	Chemical energy release
$\nu_0 = \kappa_0 = D_0$	$2.9 \times 10^{-5} \text{ g/s-cm-K}^{0.7}$	Transport constants
Output		
S_l	298 cm/s	Laminar flame speed
T_b	$7.289 T_0$	Post-flame temperature
ρ_b	$0.1372 \rho_0$	Post-flame density
x_l	0.035 cm	Laminar flame thickness
D_{CJ}	$1.993 \times 10^5 \text{ cm/s}$	CJ detonation velocity
P_{ZND}	$31.47 P_0$	Post-shock pressure
P_{CJ}	$16.24 P_0$	Pressure at CJ point
T_{ZND}	$3.457 T_0$	Post-shock temperature
T_{CJ}	$9.010 T_0$	Temperature at CJ point
ρ_{ZND}	$9.104 \rho_0$	Post-shock density
ρ_{CJ}	$1.802 \rho_0$	Density at CJ point
x_d	0.01927 cm	1D half-reaction thickness
λ	1–2 cm	Detonation cell size

A two-dimensional calculation of a propagating detonation, with enough resolution, produces reasonable detonation cell structure



**Background: 1 atm, 293 K
Stoichiometry H₂-air
Minimum dx = 1/2048 cm**

**Two-level detonation
cell structure,
 $\lambda = 1 - 2$ cm**

Hot-Spot Ignition in Stoichiometric Hydrogen-Air

Consider three different gradients.
 $\Delta x = 0.016$ cm, uniform mesh.

Hot spot created by assuming
 adiabatic mixing between hot
 burned and cold unburned gases.

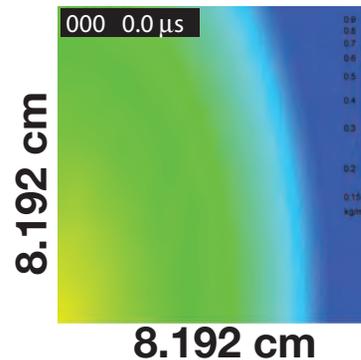
C is concentraion (mass fraction
 of unburned material)

$$\frac{dC}{dx} = \frac{0.4}{R} \quad \text{and} \quad \frac{dC}{dy} = 2 \frac{dC}{dx}$$

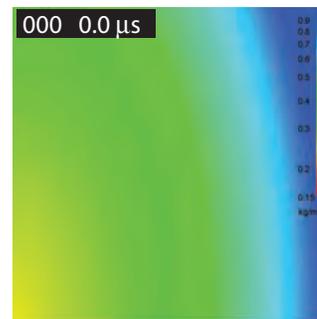
Result is an elliptical distribution.

At the center,

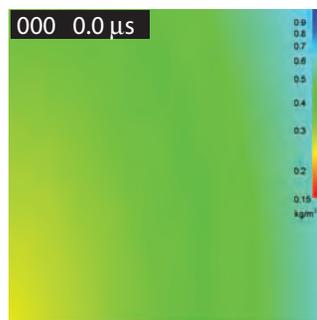
$$C = 0.6 \quad T = 1029 \text{ K}$$



$R = 7$
*Gradient too steep
 to ignite a detonation.
 Result is a decoupled
 flame and a shock*



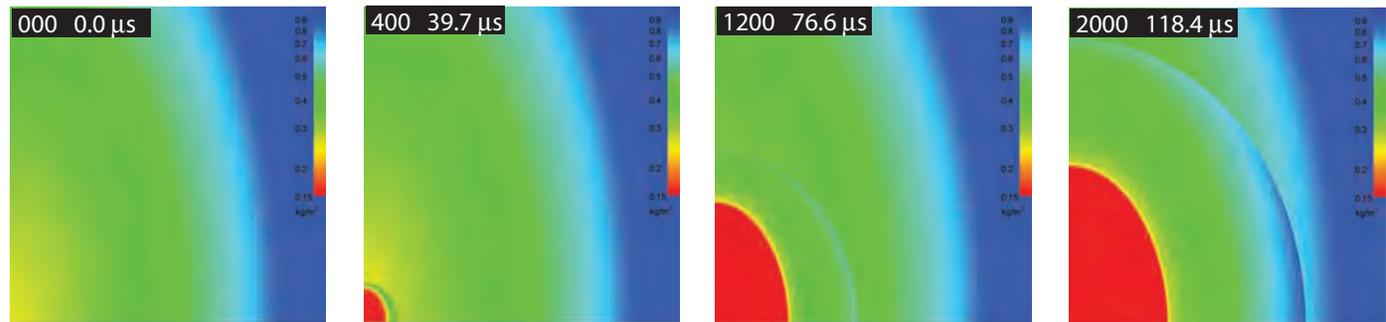
$R = 8$
 Detonation appears
 in the y-direction



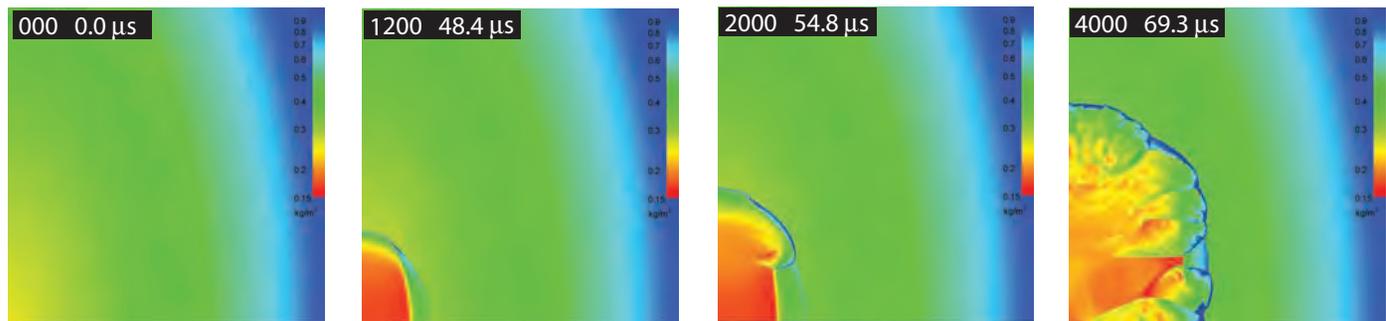
$R = 10$
 Detonation appears
 in both directions

Hot-Spot Ignition in Stoichiometric Hydrogen-Air

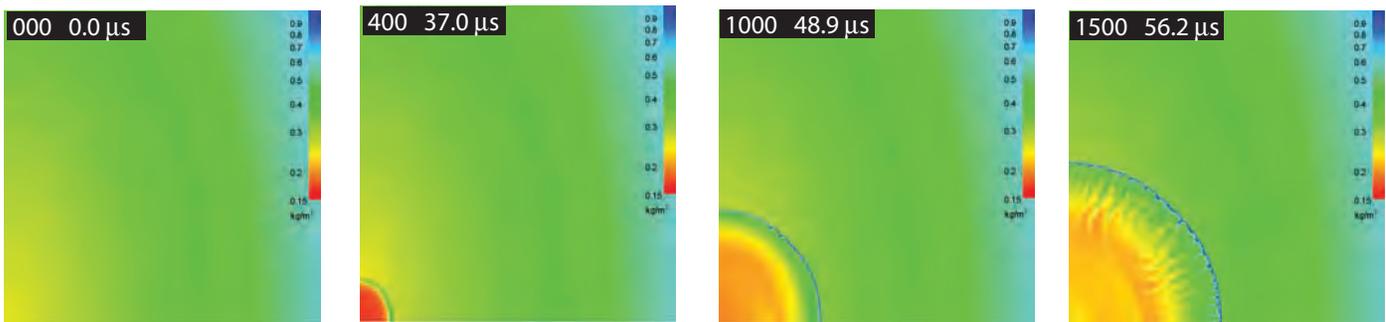
**Case 1 --
Ignition in gradient
leaves a shock
and a flame**



**Case 2 --
Ignition in gradient
- intermediate
case**



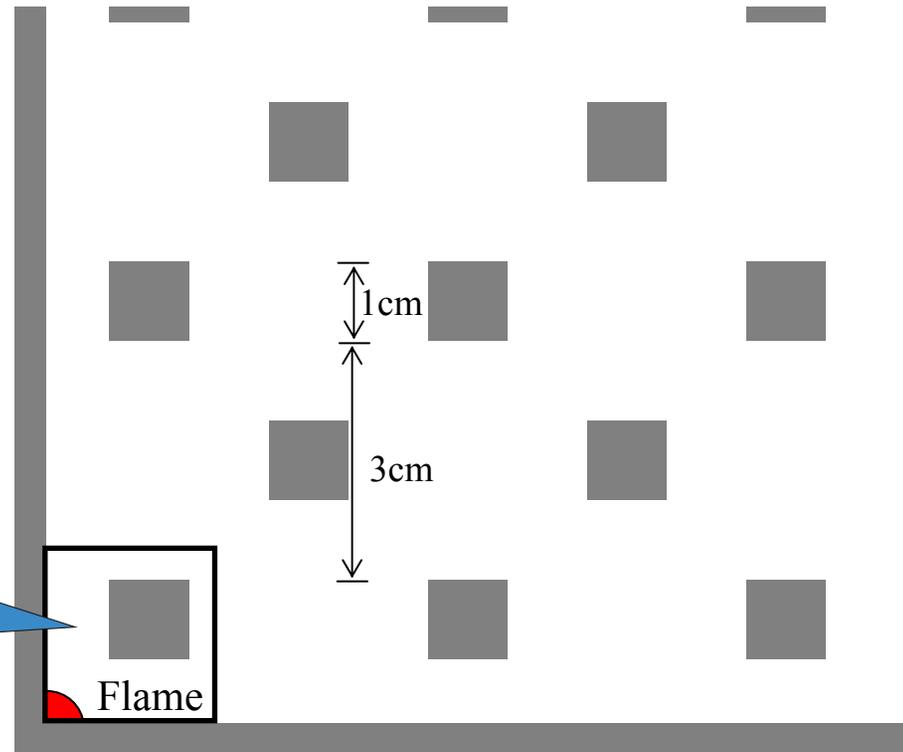
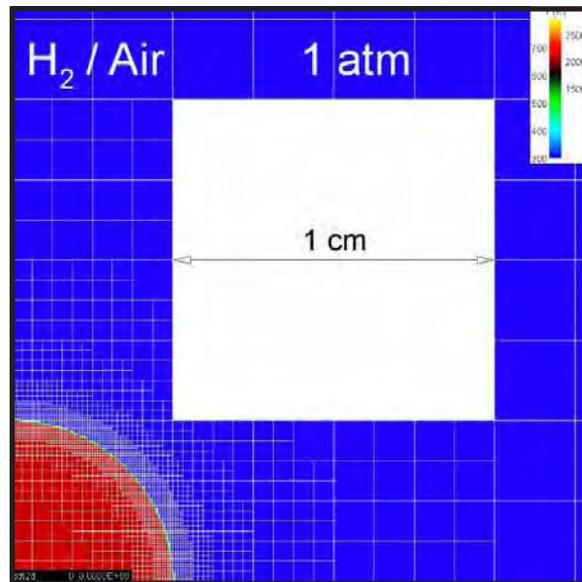
**Case 3
Ignition in gradient
produces a
detonation**



Background: unshocked:1 atm, 293 K. Each frame is 8.2 cm x 8.2 cm. dx = 0.016 cm.

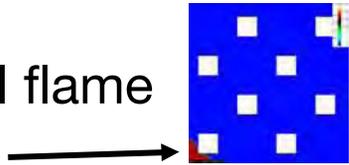
An Array of Obstacles

- Array of 1 cm x 1 cm square obstacles
- H₂-Air background
- Staggered layout

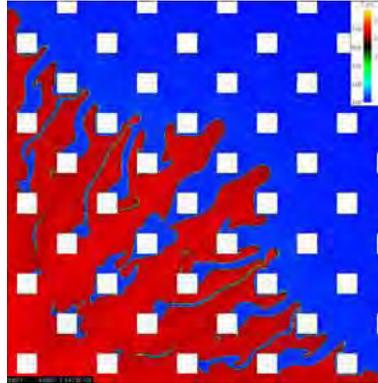


DDT in an Array of Obstacles

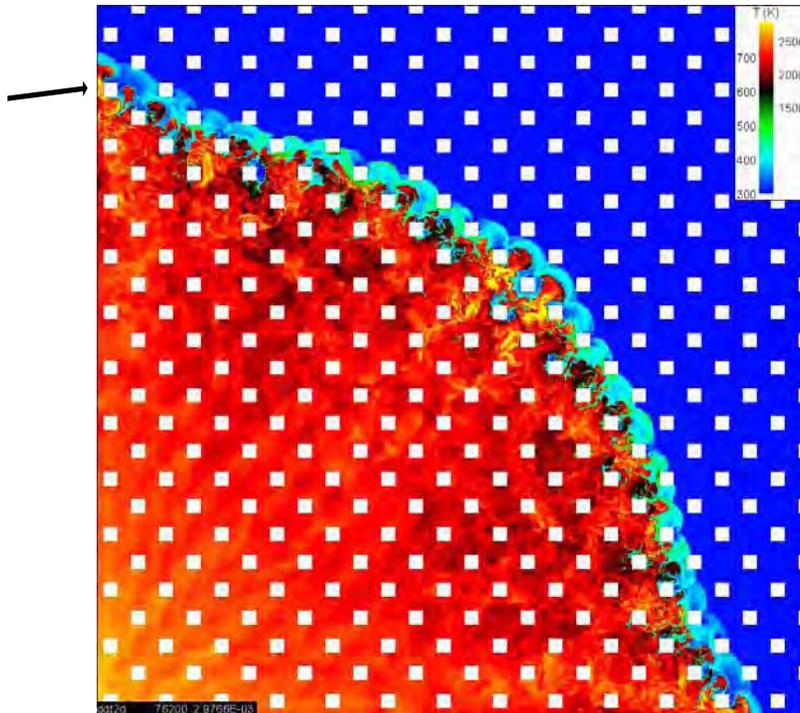
1. Initial flame



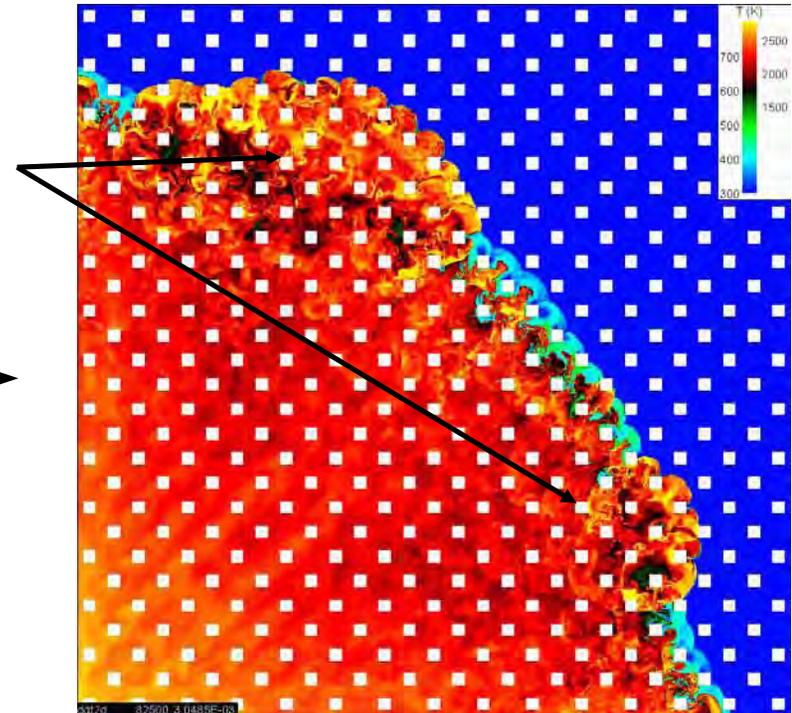
2. The flame becomes turbulent.



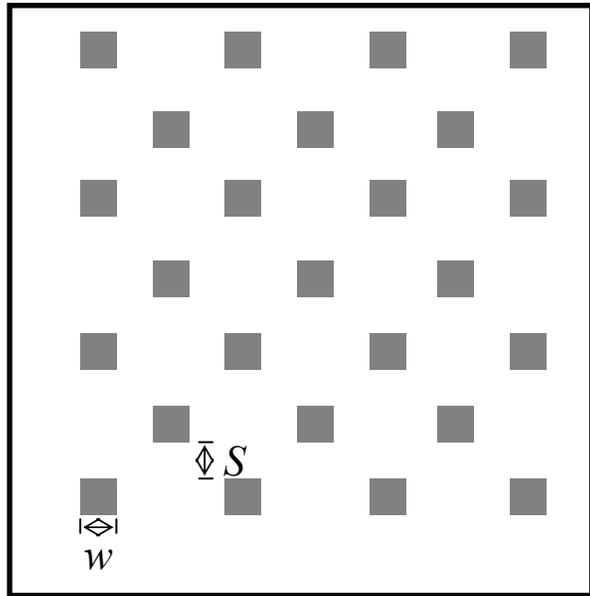
3. Shock wave formation



4. DDT: shock reflection from obstacles

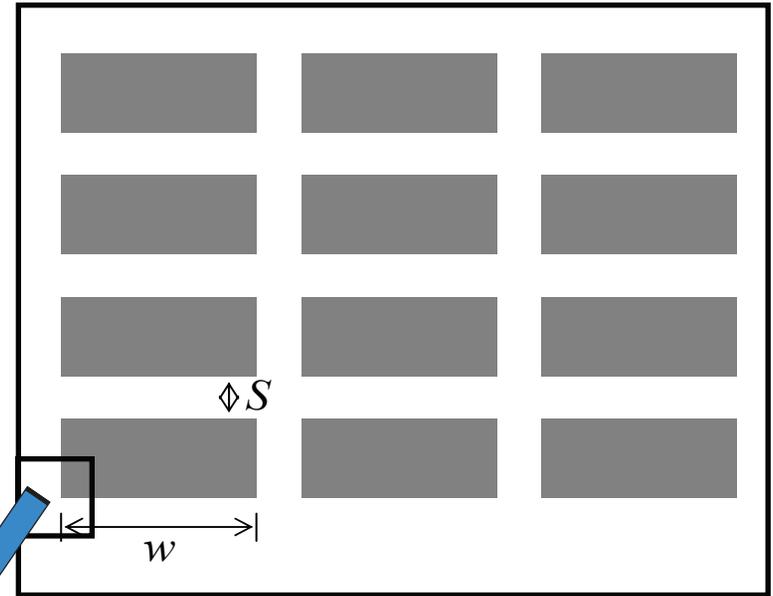


Array of Obstacles

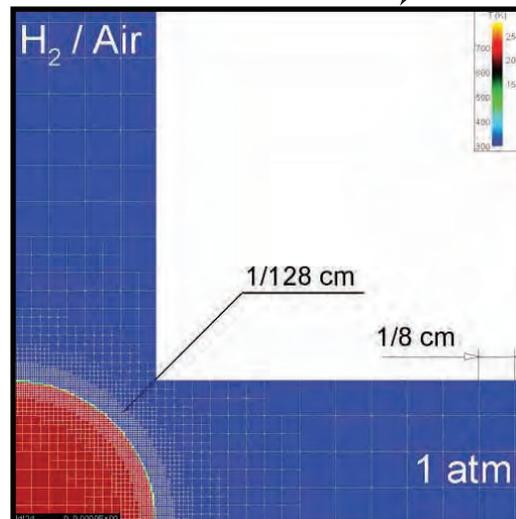


$S = 1 \text{ cm}$, $w = 1 \text{ cm}$
(staggered)

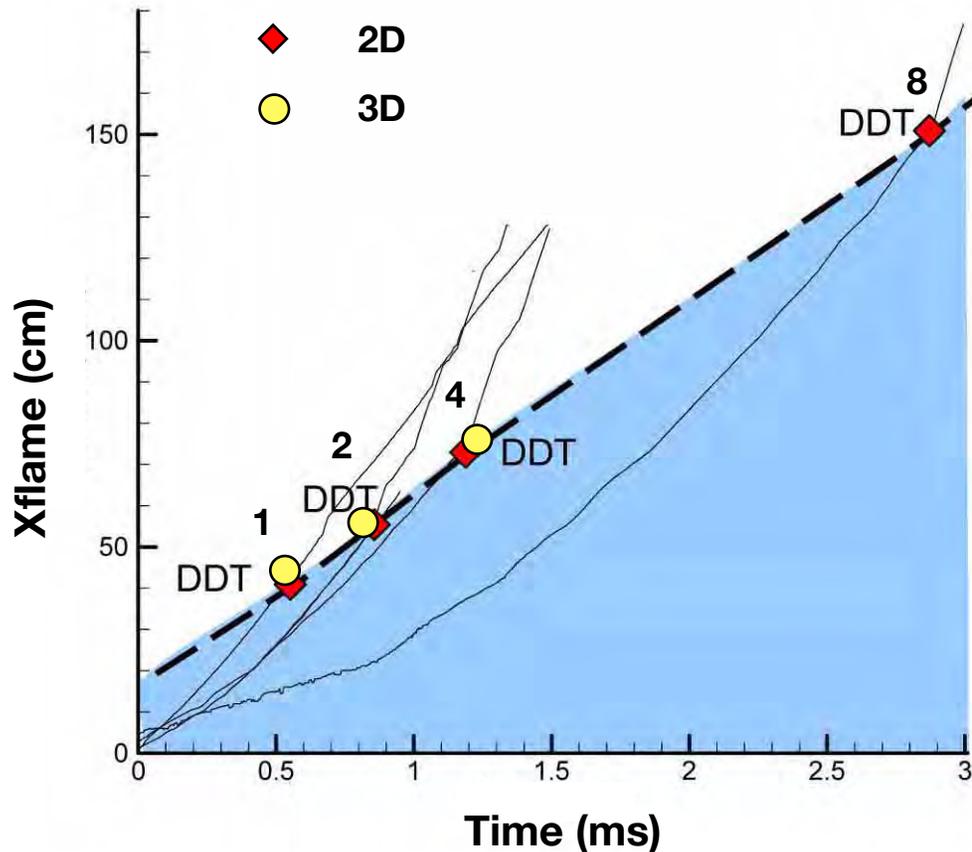
Different size and layout



$S = 1 \text{ cm}$, $w = 4 \text{ cm}$
(symmetric)



Flame Propagation and DDT in Channel with Obstacles 2D and 3D Calculations - Summary



**DDT = First successful
hot-spot transition**

**Choking flame: 1 cm
Quasi-detonation: 2, 4 cm
Detonation: 8 cm**

**Two Effects to Note:
Ignition of Hot Spots
Effect of Scale
System has to be big
enough to sustain DDT
These are related ...**

And now we have experiments (Teordorczyk, 2006)

Thoughts on the “Mysterious” Agreement ...

Richtmyer-Meshkov interactions are a source of turbulence in compressible, high-speed flows with repeated shock-flame interactions.

2D and 3D RM interactions have very similar growth rates and amplitudes -- both qualitatively and quantitatively -- in the linear regime, and differ only slightly well into the nonlinear regime.

RM can be the major source of turbulence, and when this is the case, all scales are populated simultaneously.

This gives us a clue as to why 2D and 3D simulations of high-speed deflagrations have certain similar properties.

Where do we go from here ?

Calibration with experiments

Chemical model

Effects of variable stoichiometry

Other detonable materials

Even more “difficult” geometries

Why do simulations “work” ..?

Effects of non-Kolmogorov turbulence

Stochasticity